

Advanced statistical mechanics
(WS 12/13, FU Berlin)

Problem sheet 7

Due date: December 12, 2012

Problem 21: Density matrices (1+1+1+1+1+1+1+1=8)

Let \mathcal{H} be a d -dimensional Hilbert space and denote the set of linear operators on \mathcal{H} by $\mathcal{B}(\mathcal{H}) := \{A : \mathcal{H} \rightarrow \mathcal{H} \mid A \text{ linear}\}$. A self-adjoint operator $A \in \mathcal{B}(\mathcal{H})$ is called *positive semi-definite* iff $\langle \psi | A | \psi \rangle \geq 0$ for all $|\psi\rangle \in \mathcal{H}$, which is denoted by $A \geq 0$.

a) Let $A^\dagger = A \in \mathcal{B}(\mathcal{H})$. Prove that $A \geq 0$ if and only if A has no negative eigenvalues.

The set of density matrices on \mathcal{H} is

$$\mathcal{S}(\mathcal{H}) := \{\rho \in \mathcal{B}(\mathcal{H}) : \rho \geq 0, \text{Tr}(\rho) = 1\}. \quad (1)$$

b) Let $\rho \in \mathcal{B}(\mathcal{H})$. Prove that ρ is a density matrix if and only if there is a probability vector $p \in [0, 1]^n$ and normalized state vectors $(|\psi_j\rangle)_{j=1}^n \subset \mathcal{H}$ such that

$$\rho = \sum_{j=1}^n p_j |\psi_j\rangle\langle\psi_j|. \quad (2)$$

c) Let $A^\dagger = A \in \mathcal{B}(\mathcal{H})$, i.e. A is an observable. Prove for the decomposition (2) that

$$\text{Tr}(\rho A) = \sum_{j=1}^n p_j \langle A \rangle_{\psi_j}, \quad (3)$$

i.e., that the expectation value of A in state ρ is the same as the expectation value w.r.t. p_j of the “pure state expectation values” $\langle A \rangle_{\psi_j} = \langle \psi_j | A | \psi_j \rangle$.

d) Consider the special case $\mathcal{H} = \mathbb{C}^2$. Show that the decomposition (2) is not unique, i.e., find pairwise different and normalized $|\psi_1\rangle, |\psi_2\rangle, |\phi_1\rangle, |\phi_2\rangle \in \mathbb{C}^2$ and probability vectors $p, q \in [0, 1]^2$ such that

$$\sum_{j=1}^2 p_j |\psi_j\rangle\langle\psi_j| = \sum_{j=1}^2 q_j |\phi_j\rangle\langle\phi_j|. \quad (4)$$

e) Find a necessary condition on $(|\psi_j\rangle)_{j=1}^n$ in Eq. (2) such that n is the rank of ρ .

A state $\rho \in \mathcal{S}(\mathcal{H})$ is called *pure* if there is a vector $|\psi\rangle \in \mathcal{H}$ such that $\rho = |\psi\rangle\langle\psi|$ and *mixed* if no such a vector exists.

f) Prove that $\rho \in \mathcal{S}(\mathcal{H})$ is a pure state if and only if $\text{Tr}(\rho^2) = 1$.

g) Prove that $\mathcal{S}(\mathcal{H})$ is a convex set, i.e.

$$\rho, \sigma \in \mathcal{S}(\mathcal{H}) \quad \Rightarrow \quad \lambda\rho + (1 - \lambda)\sigma \in \mathcal{S}(\mathcal{H}) \quad \forall \lambda \in [0, 1]. \quad (5)$$

Remark: This implies that a probability distribution over quantum states naturally defines again a quantum state.

A point $\rho \in \mathcal{S}$ of a convex set \mathcal{S} is called an *extremal point* of \mathcal{S} if it is impossible to write it as a non-trivial convex combination of two other points, i.e., for all $\rho_1, \rho_2 \in \mathcal{S}$ and $\lambda \in]0, 1[$

$$\rho = \lambda\rho_1 + (1 - \lambda)\rho_2 \quad \Rightarrow \quad \rho_1 = \rho_2. \quad (6)$$

h) Prove that the extremal points of $\mathcal{S}(\mathcal{H})$ are exactly the pure states.

Hint: First, conclude that the pure state are exactly the ones of rank one.

Problem 22: Tensor product (1+1+1+1+2)

Let \mathcal{H} be a d -dimensional Hilbert space with basis $\{|j\rangle\}_{j=1}^d$. The dual space \mathcal{H}^* has the basis $\{\langle j|\}_{j=1}^d$, satisfying the orthonormality relation $\langle j|k\rangle = \delta_{j,k}$.

Similarly, let \mathcal{H}_i with $i = 1, 2$ be two other d_i -dimensional Hilbert space with basis $\{|j\rangle_i\}_{j=1}^{d_i}$. We will often use the common notation

$$|j, k\rangle := |j\rangle |k\rangle := |j\rangle \otimes |k\rangle := |j\rangle_1 \otimes |k\rangle_2. \quad (7)$$

The vectors $\{|j, k\rangle\}_{j=1, \dots, d_1; k=1, \dots, d_2}$ are an orthonormal basis of the (tensor) product space $\mathcal{H}_1 \otimes \mathcal{H}_2$ and satisfy the orthogonality relations

$$\langle j, k|l, m\rangle = \langle j|l\rangle \langle k|m\rangle = \delta_{j,l} \delta_{k,m}. \quad (8)$$

The tensor product of two arbitrary vectors $|\psi\rangle \in \mathcal{H}_1$ and $|\phi\rangle \in \mathcal{H}_2$ can be defined via bilinear extension as

$$|\psi\rangle \otimes |\phi\rangle := \sum_{j=1}^{d_1} \sum_{k=1}^{d_2} \langle j|\psi\rangle \langle k|\phi\rangle |j, k\rangle. \quad (9)$$

- What is the Hilbert space dimension of the n -fold tensor product $\mathcal{H}^{\otimes n} = \mathcal{H} \otimes \dots \otimes \mathcal{H}$?
- For $|\psi\rangle \in \mathcal{H}_1$ and $|\phi\rangle, |\xi\rangle \in \mathcal{H}_2$ show that

$$|\psi\rangle \otimes |\phi\rangle + |\psi\rangle \otimes |\xi\rangle = |\psi\rangle (|\phi\rangle + |\xi\rangle). \quad (10)$$

There are the natural isomorphisms

$$\mathcal{H} \otimes \mathcal{H}^* \rightarrow \mathcal{B}(\mathcal{H}), \quad |j\rangle \otimes \langle k| \mapsto |j\rangle \langle k|, \quad (11)$$

$$\mathcal{H}_1 \otimes \mathcal{H}_2 \rightarrow \mathcal{H}_2 \otimes \mathcal{H}_1, \quad |j\rangle \otimes |k\rangle \mapsto |k\rangle \otimes |j\rangle \quad (12)$$

and consequently

$$\mathcal{B}(\mathcal{H}_1) \otimes \mathcal{B}(\mathcal{H}_2) \rightarrow \mathcal{B}(\mathcal{H}_1 \otimes \mathcal{H}_2), \quad |j\rangle \langle k| \otimes |l\rangle \langle m| \mapsto |j, l\rangle \langle k, m|. \quad (13)$$

Therefore, one often writes, e.g., $|j\rangle \langle k| \otimes |l\rangle \langle m| = |j, l\rangle \langle k, m|$. Importantly, $|\psi\rangle \otimes |\phi\rangle \neq |\phi\rangle \otimes |\psi\rangle$ in general, even if $\mathcal{H}_1 \otimes \mathcal{H}_2 \cong \mathcal{H}_2 \otimes \mathcal{H}_1$. So one needs to be careful here: If one reorders basis elements during one calculation one needs to reorder *all* basis elements in the same way!

In terms of the basis (11) an operator $A \in \mathcal{B}(\mathcal{H})$ can be written as

$$A = \sum_{j,k=1}^d a_{j,k} |j\rangle \langle k|. \quad (14)$$

- How can one calculate the $a_{j,k}$?
- For the special case, where $A = \sigma_x \otimes \sigma_z$ find the basis expansion (14) of A explicitly, where σ_x, σ_y , and σ_z denote the Pauli matrices.
- Show that for $A, C \in \mathcal{B}(\mathcal{H}_1)$ and $B, D \in \mathcal{B}(\mathcal{H}_2)$

$$(A \otimes B)(C \otimes D) = (AC) \otimes (BD). \quad (15)$$

Problem 23: Two indistinguishable particles (1+1+1+1+1)

- a) Consider two state vectors $|\psi\rangle$ and $|\phi\rangle$ that represent indistinguishable states of a quantum system. Which mathematical relation between $|\psi\rangle$ and $|\phi\rangle$ is necessary and sufficient for indistinguishability? Explain your answer!

The swap operator $S \in \mathcal{B}(\mathcal{H} \otimes \mathcal{H})$ is defined via

$$S |j\rangle |k\rangle = |k\rangle |j\rangle, \quad (16)$$

where $\{|j\rangle\}_j$ is some orthonormal basis of the Hilbert space \mathcal{H} .

- b) Show that the eigenvalues of S are $\text{spec}(S) = \{-1, 1\}$.

From now on we consider the case, where the one-particle Hilbert space is two-dimensional, i.e., $\mathcal{H} \cong \mathbb{C}^2$.

- c) Show that the famous *Bell states*

$$|\phi^+\rangle = (|0, 0\rangle + |1, 1\rangle) / \sqrt{2}, \quad |\phi^-\rangle = (|0, 0\rangle - |1, 1\rangle) / \sqrt{2}, \quad (17)$$

$$|\psi^+\rangle = (|0, 1\rangle + |1, 0\rangle) / \sqrt{2}, \quad |\psi^-\rangle = (|0, 1\rangle - |1, 0\rangle) / \sqrt{2} \quad (18)$$

are an orthonormal basis for our two-particle Hilbert-space.

- d) For each eigenspace of S find a basis in terms of Bell states.
- e) Now consider quantum state $|\psi\rangle$ of two indistinguishable particles, one of which is in state $|\psi_1\rangle \in \mathbb{C}^2$ and the other in state $|\psi_2\rangle \in \mathbb{C}^2$. Conclude that $|\psi\rangle$ is either symmetric or anti-symmetric under exchange of the two particles.

Remark: If $|\psi\rangle$ is symmetric the particles are called *bosons* and if $|\psi\rangle$ is anti-symmetric then they are called *fermions*.